

# Preconditioning for Nonsymmetry and Time-dependence

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joint work with  
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# Iterative methods

For self-adjoint problems/symmetric matrices, iterative methods of choice exist: conjugate gradients for SPD, MINRES otherwise

but many possible methods for non-self-adjoint problems/nonsymmetric matrices: GMRES , BICGSTAB , QMR , IDR , . . .

# Iterative methods

Arises because of convergence guarantees:

- for symmetric matrices: descriptive convergence bounds  $\Rightarrow$  a priori estimates of iterations for acceptable convergence; good preconditioning ensures fast convergence.
- for nonsymmetric matrices: by contrast, to date there are no generally applicable *and descriptive* convergence bounds even for GMRES ; for any of the other nonsymmetric methods without a minimisation property, convergence theory is extremely limited  $\Rightarrow$  no good a priori way to identify what are the desired qualities of a preconditioner

A major theoretical difficulty, but heuristic ideas abound!

The situation is more severe than this:

Theorem (*Greenbaum, Ptak and Strakos, 1996*)

Given any set of eigenvalues and any monotonic convergence curve, then for any  $b$  there exists a matrix  $B$  having those eigenvalues and an initial guess  $x_0$  such that GMRES for  $Bx = b$  with  $x_0$  as starting vector will give that convergence curve.

In fact more extreme negative results than this exist.

One way to address such questions:  
look for (non-standard) inner products in which a problem  
might be self-adjoint

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might be self-adjoint

such inner products exist for a real nonsymmetric matrix  $B$   
if and only if  $B$  is diagonalizable and has real eigenvalues

but preconditioners still would have to have self-adjointness  
in any relevant non-standard inner product!!

# Convection-diffusion equation:

$$\frac{\partial u}{\partial t} + \mathbf{b} \cdot \nabla u - \epsilon \nabla^2 u = f \quad \text{in } \Omega \times (0, T], \quad \Omega \subset \mathbb{R}^2 \text{ or } \mathbb{R}^3$$

$$u(\mathbf{x}, 0) = u_0(\mathbf{x}), \quad u \text{ given on } \partial\Omega$$

- arises widely e.g. as a subproblem in Navier-Stokes
- is *non-self-adjoint*  $\Rightarrow$  nonsymmetric discretization matrix

$\Rightarrow$  convergence of Krylov subspace methods not easily described

so no mathematical idea how to precondition

For steady convection-diffusion

$$\mathbf{b} \cdot \nabla u - \epsilon \nabla^2 u = f$$

the nonsymmetric issue remains even in 1-dimension

$$u' - \epsilon u'' = f$$

(though iterative methods not so crucial here!)



$$u' - \epsilon u'' = f$$

is

$$\left(-\exp(-x/\epsilon)u'\right)' = \frac{1}{\epsilon}\exp(-x/\epsilon)f$$

so continuous problem is self-adjoint in a highly distorted inner product based on this integrating factor (given certain simple boundary conditions)

Discretizations however have matrices which are *not* self-adjoint in any inner product for large enough  $h/\epsilon$

Recently (*Pestana & W, 2015*) we have made progress for real nonsymmetric Toeplitz (constant diagonal) matrices regardless of nonnormality with a very simple observation: If  $B$  is a real Toeplitz matrix then

$$\underbrace{\begin{bmatrix} a_0 & a_{-1} & \cdot & \cdot & a_{1-n} \\ a_1 & a_0 & a_{-1} & \cdot & \cdot \\ \cdot & a_1 & a_0 & \cdot & \cdot \\ \cdot & \cdot & \cdot & \cdot & a_{-1} \\ a_{n-1} & \cdot & \cdot & a_1 & a_0 \end{bmatrix}}_B \quad \underbrace{\begin{bmatrix} 0 & 0 & \cdot & 0 & 1 \\ 0 & \cdot & 0 & 1 & 0 \\ \cdot & 0 & 1 & 0 & \cdot \\ \cdot & \cdot & \cdot & \cdot & \cdot \\ 1 & 0 & \cdot & 0 & 0 \end{bmatrix}}_Y$$

is the real *symmetric* matrix

$$\begin{bmatrix} a_{1-n} & \cdot & \cdot & a_{-1} & a_0 \\ \cdot & \cdot & a_{-1} & a_0 & a_1 \\ \cdot & \cdot & a_0 & a_1 & \cdot \\ a_{-1} & \cdot & \cdot & \cdot & \cdot \\ a_0 & a_1 & \cdot & \cdot & a_{n-1} \end{bmatrix}$$

Thus MINRES can be robustly applied to  $BY$  — it is symmetric but generally indefinite — and its convergence will depend only on eigenvalues.

BUT preconditioning? – needs to be symmetric and positive definite for MINRES

Fortunately it is well known that Toeplitz matrices are well preconditioned by related circulant matrices,  $C$  (Strang, 1986, Chan, 1988) which are diagonalised by an FFT in  $O(n \log n)$  work:  $C = U^* \Lambda U$ .

For many Toeplitz matrices we have that the Strang or Optimal (Chan) circulant  $C$  satisfy

$$C^{-1}B = I + R + E$$

where  $R$  is of small rank and  $E$  is of small norm

⇒ eigenvalues clustered around 1 except for a few outliers

To ensure symmetric and positive definite preconditioner for  $BY$  just use the circulant matrix

$$|C| = U^* |\Lambda| U$$

which is real symmetric and positive definite

Theorem (*Pestana & W, 2015*)

$$|C|^{-1} BY = J + R + E$$

where  $J$  is real symmetric and orthogonal with eigenvalues  $\pm 1$ ,  $R$  is of small rank and  $E$  is of small norm

$\Rightarrow$  guaranteed fast convergence because MINRES convergence only depends on eigenvalues which are clustered around  $\pm 1$  except for few outliers!

# Example

$$B = \begin{bmatrix} 1 & 0.01 & & & & \\ 1 & 1 & 0.01 & & & \\ & \ddots & \ddots & \ddots & & \\ & & 1 & 1 & 0.01 & \\ & & & 1 & 1 & \\ & & & & 1 & 1 \end{bmatrix} \in \mathbb{R}^{n \times n}$$

$n$	Condition number of $B$	preconditioned MINRES iters
10	14	6
100	207	6
1000	$2.6 \times 10^6$	6

Similar ideas apply for block Toeplitz matrices—higher dimensions—but theory not so good

# Time-dependent Convection-Diffusion

$$\frac{\partial u}{\partial t} + \mathbf{b} \cdot \nabla u - \epsilon \nabla^2 u = f \quad \text{in } \Omega \times (0, T], \quad \Omega \subset \mathbb{R}^2 \text{ or } \mathbb{R}^3$$

$$u(\mathbf{x}, 0) = u_0(\mathbf{x}), \quad u \text{ given on } \partial\Omega$$

Finite elements in space ( $\mathbf{x}$ ),  $\theta$  time stepping gives

$$M \frac{u_k - u_{k-1}}{\tau} + K \left( \theta u_k + (1 - \theta) u_{k-1} \right) = \mathbf{f}_k$$

$M \in \mathbb{R}^{n \times n}$ : SPD mass matrix (identity operator, but same sparsity as  $K$ )

$K \in \mathbb{R}^{n \times n}$ : nonsymmetric discrete steady convection-diffusion matrix

Rearranging:

$$\left( M + \tau \theta K \right) \mathbf{u}_k = \left( M - \tau (1 - \theta) K \right) \mathbf{u}_{k-1} + \tau \mathbf{f}_k,$$

$$k = 1, 2, \dots, N$$

$$N\tau = T$$

Recall for unconditional stability:  $\frac{1}{2} \leq \theta \leq 1$

$\theta = 1$ : backwards Euler,       $\theta = \frac{1}{2}$ : Crank-Nicolson

else need  $\tau = \mathcal{O}(h^2)$ : very small time steps for explicit method

$$\left( M + \tau \theta K \right) \mathbf{u}_k = \left( M - \tau (1 - \theta) K \right) \mathbf{u}_{k-1} + \tau \mathbf{f}_k,$$

$$k = 1, 2, \dots, N$$

Standard solution method:

Solve the  $N$  separate  $n \times n$  nonsymmetric linear systems (sequentially) for  $k = 1, 2, \dots, N$ . Here we use a geometric multigrid due to Ramage specifically for convection diffusion

$\Rightarrow r = 5$  V-cycles for solution of each linear system to a relative residual tolerance of  $10^{-6}$

Hence if we (quite reasonably) regard 1 V-cycle as the main unit of work

$\Rightarrow Nr$  V-cycles sequentially for the overall solution



## Alternatively:

Write all timesteps at one go (all-at-once method):

$$\mathcal{A} \begin{bmatrix} \mathbf{u}_1 \\ \mathbf{u}_2 \\ \vdots \\ \mathbf{u}_N \end{bmatrix} = r.h.s$$

where  $\mathcal{A}$  is the matrix

$$\begin{bmatrix} M+\tau\theta K & 0 & 0 & 0 \\ -M+\tau(1-\theta)K & M+\tau\theta K & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & -M+\tau(1-\theta)K & M+\tau\theta K \end{bmatrix}$$

and  $r.h.s. = [M-\tau(1-\theta)K \mathbf{u}_0 + \tau \mathbf{f}_1, \tau \mathbf{f}_2, \dots, \tau \mathbf{f}_N]^T$

nonsymmetric!

$$\mathcal{A} = \begin{bmatrix} M + \tau\theta K & 0 & 0 & 0 \\ -M + \tau(1 - \theta)K & M + \tau\theta K & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & -M + \tau(1 - \theta)K & M + \tau\theta K \end{bmatrix}$$

$$\mathcal{A} \in \mathbb{R}^{L \times L}, \quad L = Nn$$

We propose to solve this huge linear system (for the solution at all time steps) by GMRES (or BICGSTAB) with block diagonal preconditioner

$$\mathcal{P} = \begin{bmatrix} (M + \tau\theta K)_{MG} & 0 & 0 & 0 \\ 0 & (M + \tau\theta K)_{MG} & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & 0 & (M + \tau\theta K)_{MG} \end{bmatrix}$$

where  $(M + \tau\theta K)_{MG}$  is one GMG V-cycle exactly as above.

Any other approximate nonsymmetric solver could be used.

Theory: If we used

$$\mathcal{P}_{\text{exact}} = \begin{bmatrix} (M+\tau\theta K) & 0 & 0 & 0 \\ 0 & (M+\tau\theta K) & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & 0 & (M+\tau\theta K) \end{bmatrix}$$

as preconditioner (no multigrid approximation) then we would have

$$\mathcal{P}_{\text{exact}}^{-1} \mathcal{A} = \begin{bmatrix} I & 0 & 0 & 0 \\ J & I & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & J & I \end{bmatrix},$$

$$J = (M+\tau\theta K)^{-1} (-M+\tau(1-\theta)K)$$

For

$$\mathcal{P}_{\text{exact}}^{-1} \mathcal{A} = \begin{bmatrix} I & 0 & 0 & 0 \\ J & I & 0 & 0 \\ 0 & \ddots & \ddots & 0 \\ 0 & 0 & J & I \end{bmatrix},$$

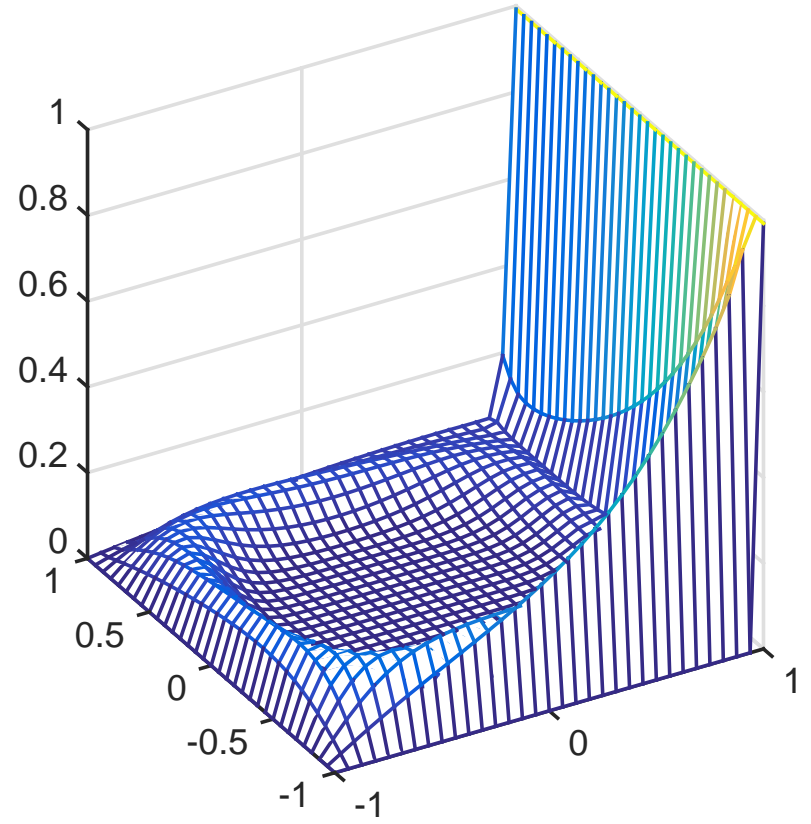
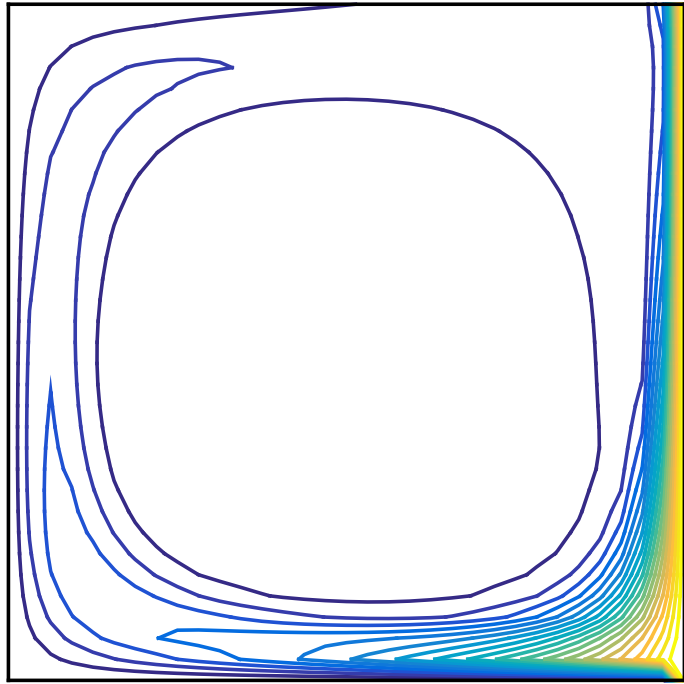
the minimum polynomial is  $(1 - s)^N$ , so GMRES would terminate (in exact arithmetic) in  $N$  iterations

We observe that  $(M + \tau\theta K)_{MG}$  is close to  $(M + \tau\theta K)$  so that convergence to a tolerance much less than the discretization error is achieved in  $N$  iterations also with  $\mathcal{P}$  as preconditioner.

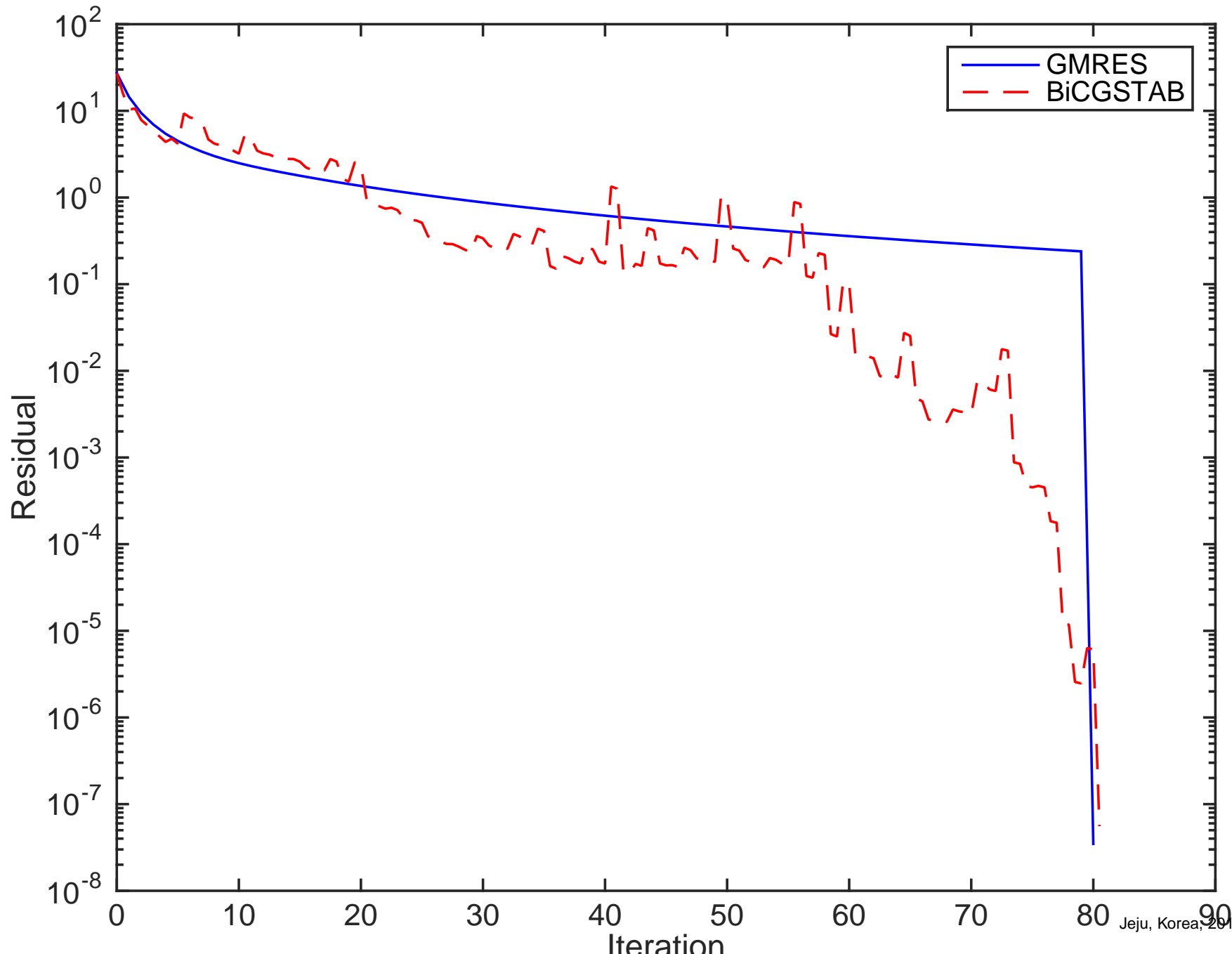
Thus:  $N$  V-cycles for each of  $N$  GMRES iterations—hence  $N^2$  ( $> Nr$ ) overall.

but with  $N$  processors, solution with  $\mathcal{P}$  is (embarrassingly) parallel—block diagonal  $\Rightarrow$  independent computation.

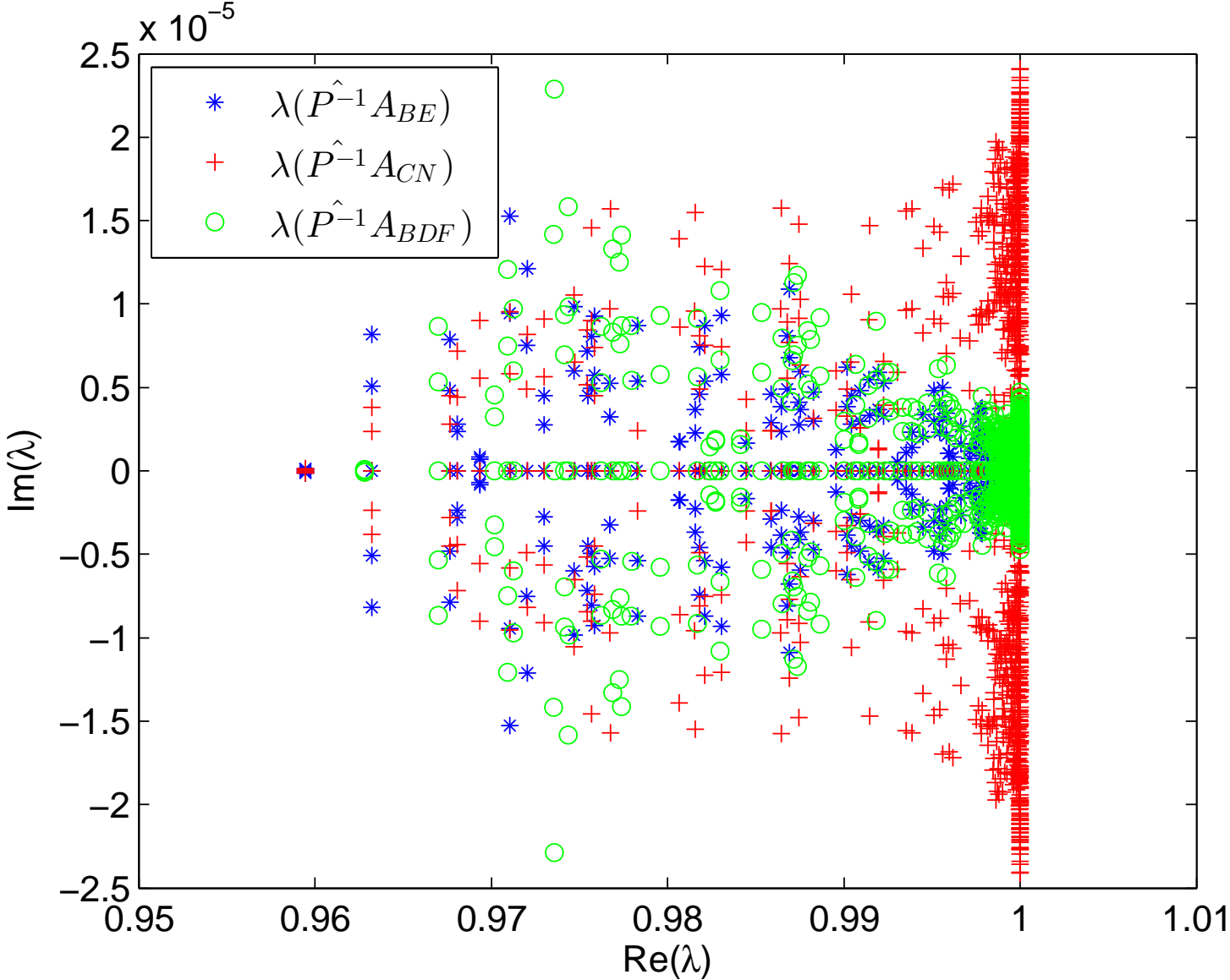
Thus parallel effort is  $N < Nr$  (= sequential effort).



# Convection-diffusion



# Convection-diffusion





# References and Acknowledgement

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